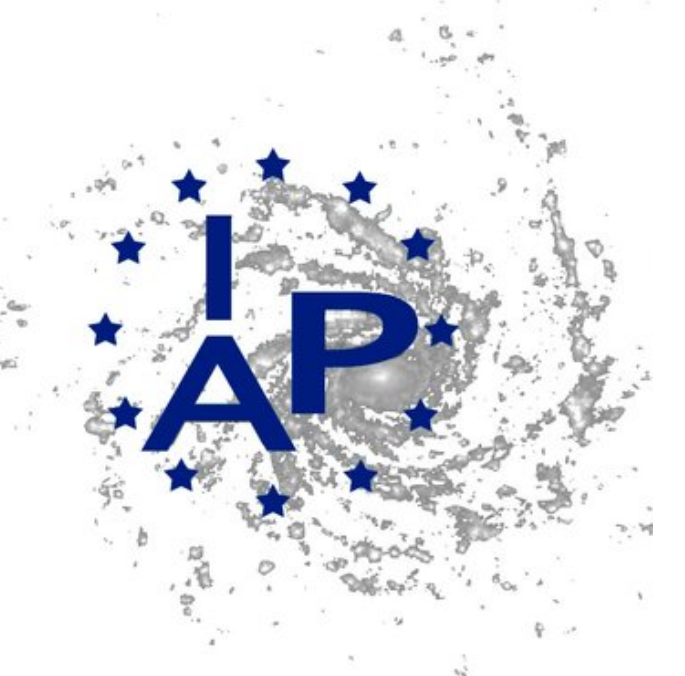




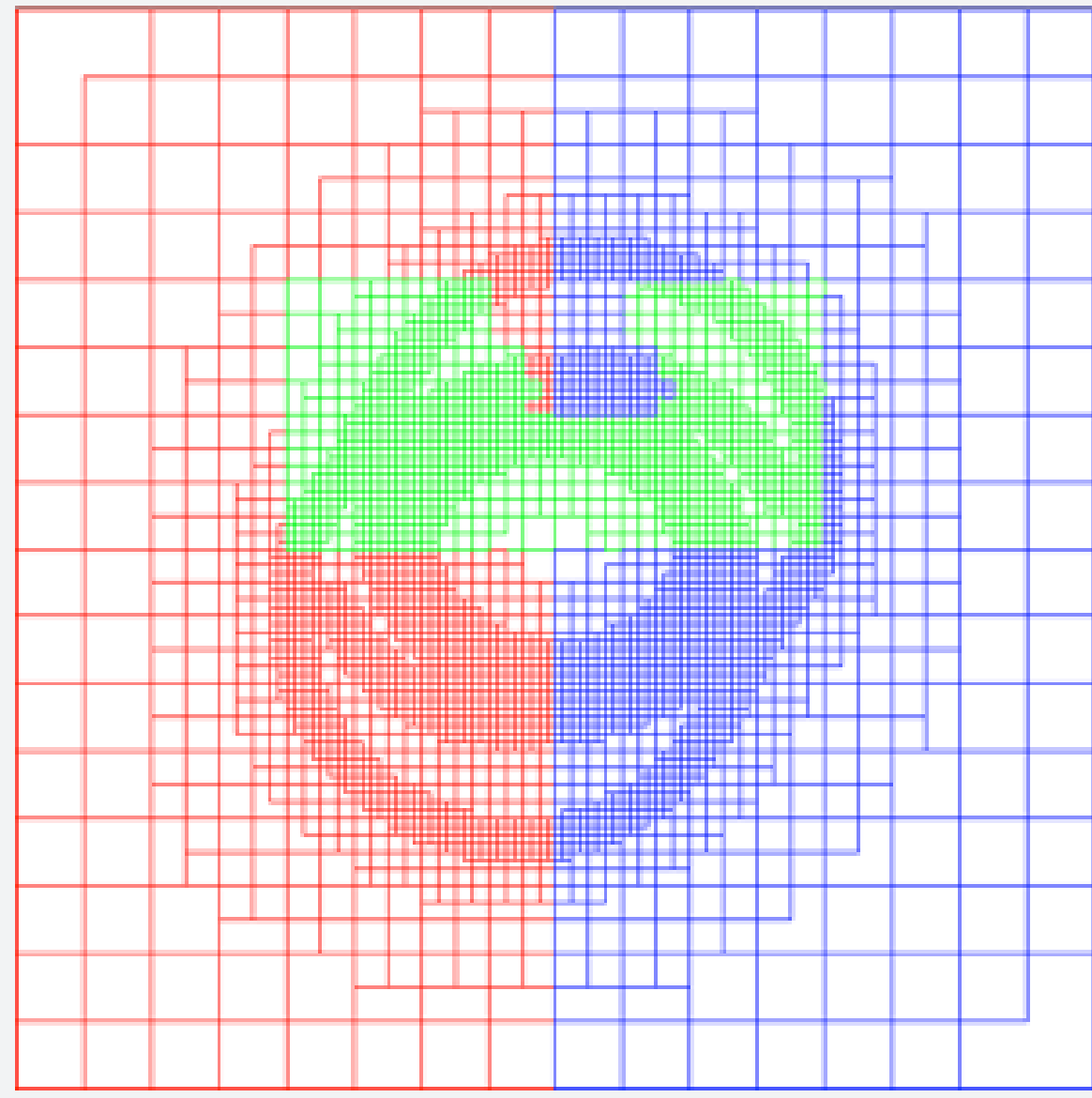
Monte Carlo tracers implemented for the adaptive mesh refinement code RAMSES



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NEED FOR LAGRANGIAN HISTORY

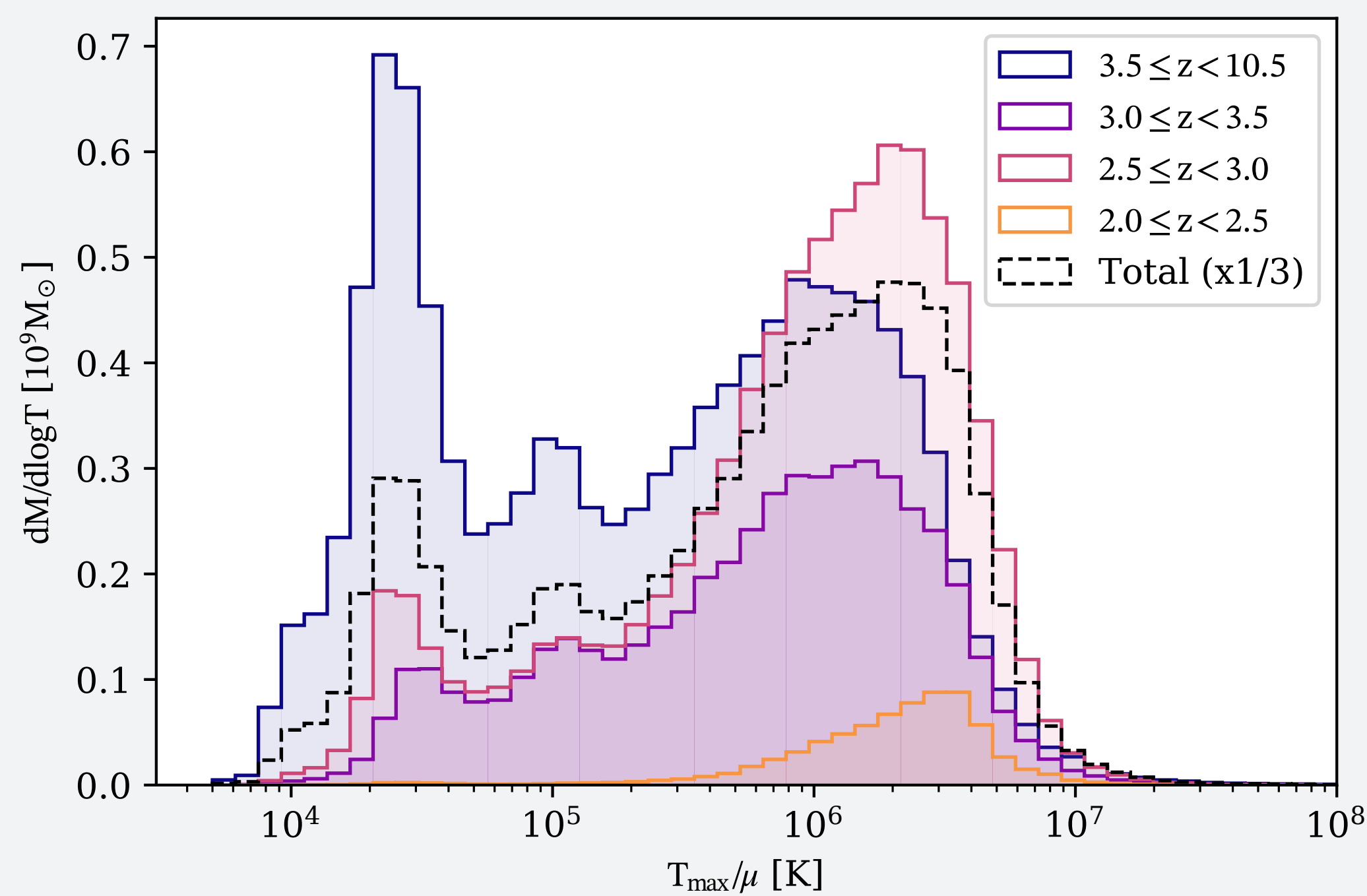
Numerical hydrodynamical simulations solving the Euler equations usually use one of two approaches. In **Smooth Particle Hydrodynamics** (SPH, e.g. [1]), the equations are solved in a Lagrangian framework. Other codes use an Eulerian approach and solve the equations on a grid: they are called **Adaptive Mesh Refinement** codes (AMR codes, e.g. [2]). AMR codes have a very flexible approach to tweak the spatial resolution. Yet they fail at providing the **Lagrangian history** of the gas.



To tackle this issue, one needs to add tracer particles in the code. The particles should have the same spatial distribution as the gas.

SCIENCE CASE

When studying the accretion of gas at high redshift $z \lesssim 2$, [3] found that in SPH simulations, most of the accreted gas has never heated above $T_{\max} \sim 10^5$ K. To study these processes in AMR code, one needs reliable tracer particles to follow the evolution of a parcel of gas as it falls onto a galaxy, eventually shock-heats at the virial radius.



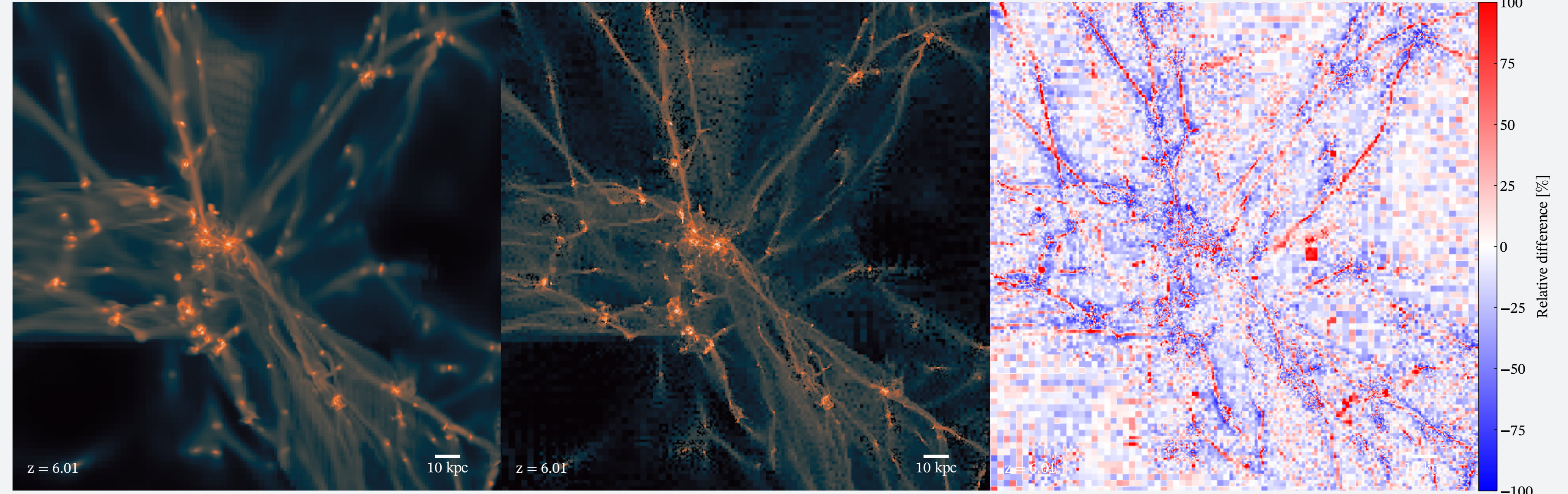
With Lagrangian information we recover their result: the accretion is mostly bimodal with anisotropic, cold neutral gas at $T_{\max} \sim 2 \times 10^4$ K and diffuse, shock-heated gas at $T_{\max} \sim 10^6$ K. The advantage of AMR codes is that they can trigger super-Lagrangian refinement, with, e.g., refinement criteria based on shock structure of the gas, vorticity of the flow, etc., with tracer particles still providing the Lagrangian history! This can be used to e.g. resolve the turbulence in cosmic filaments and track how their structure evolves as they fall onto galaxies.

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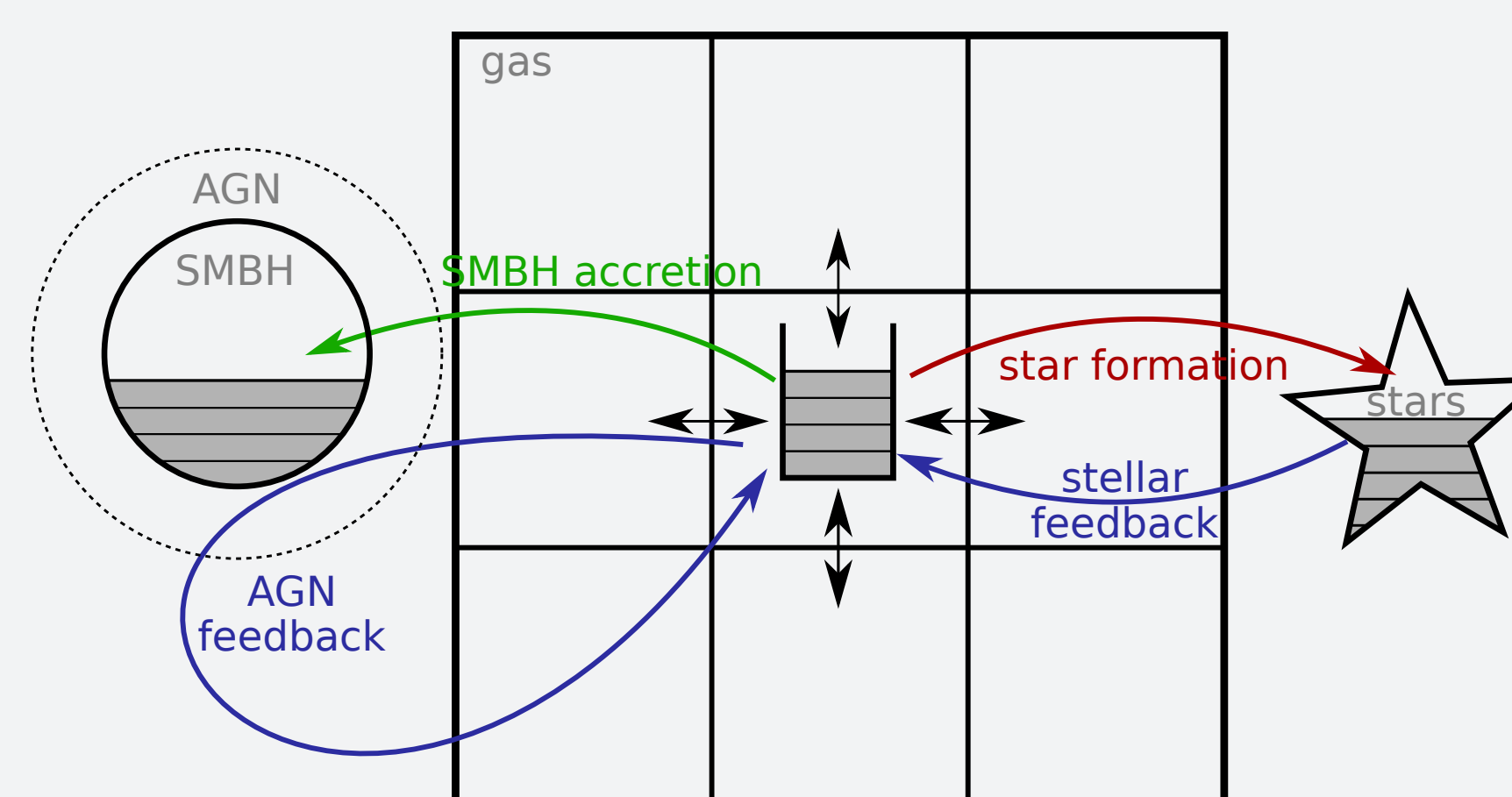
VELOCITY TRACERS

The particles are moved based on the *interpolated local velocity* v_i . The position is updated using $x_i^{t+1} = x_i^t + v_i^t \Delta t$.



Problem: the velocity-advected tracers have a distribution that is **overdense in convergent regions** (filaments, nodes) and **underdense in diverging regions** (voids), see e.g. [4]. This is due to the divergent term not being taken into account ($\rho \nabla \cdot v$).

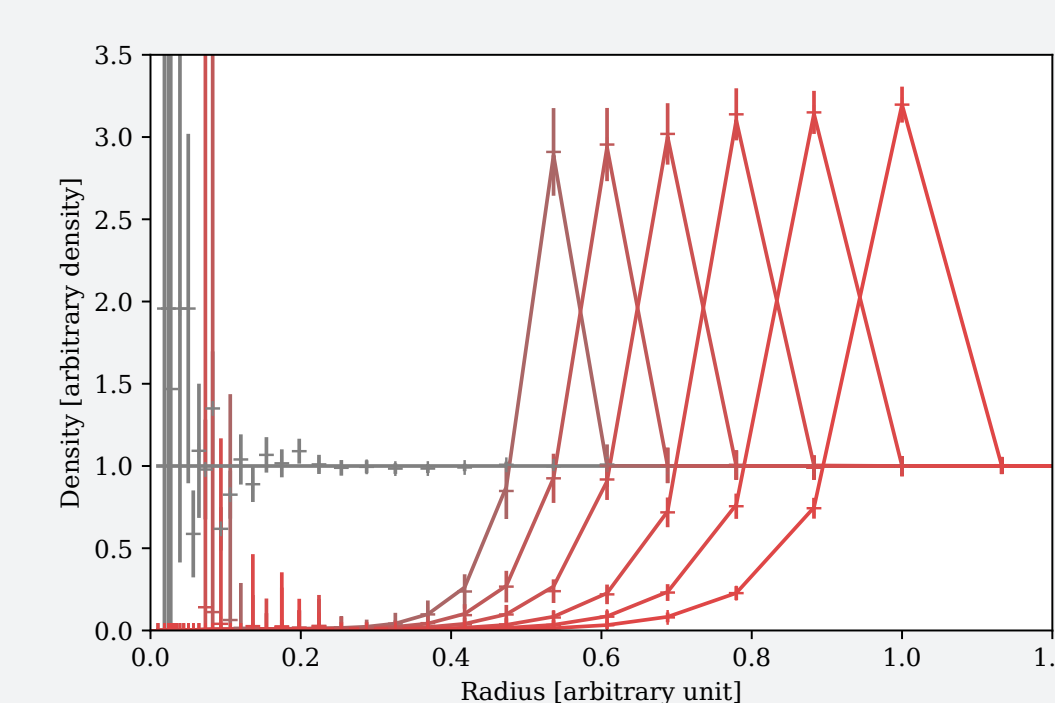
MONTE CARLO TRACERS – PRINCIPLE



The particles are moved stochastically so that the **tracer flux** equals the **gas flux** (see [5]). For a particle in cell i , the probability of jumping in cell j if the mass transfer is $\Delta M_{i \rightarrow j}$ reads

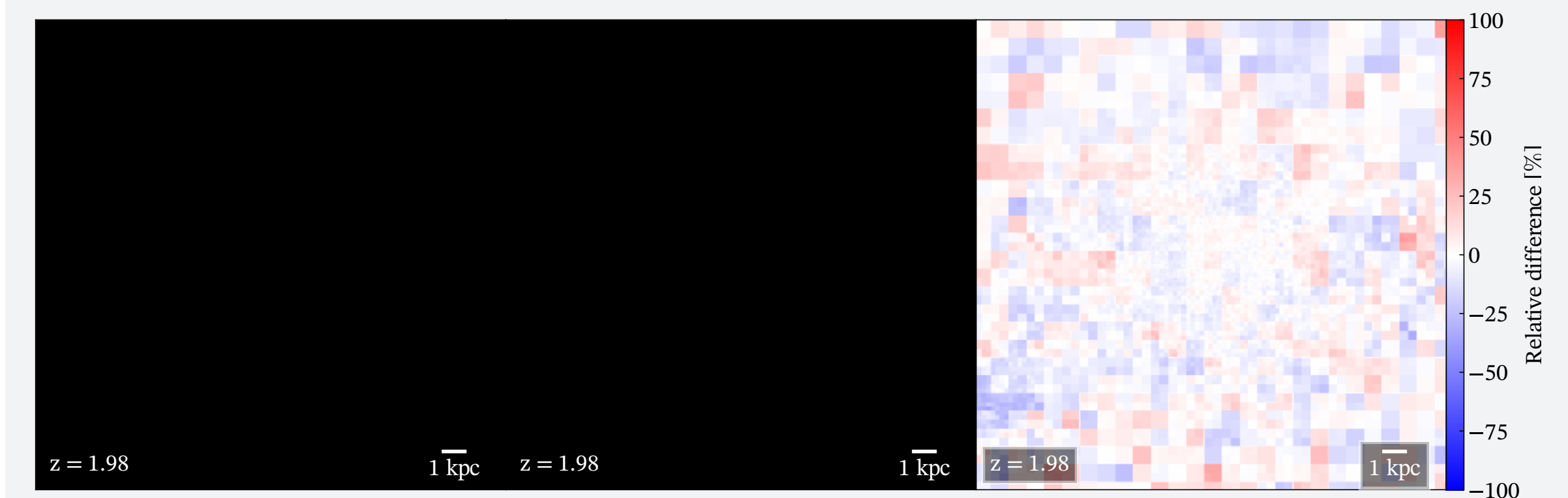
$$p_{i \rightarrow j} \equiv \frac{\Delta M_{i \rightarrow j}}{M_i}. \quad (1)$$

SEDOV TEST



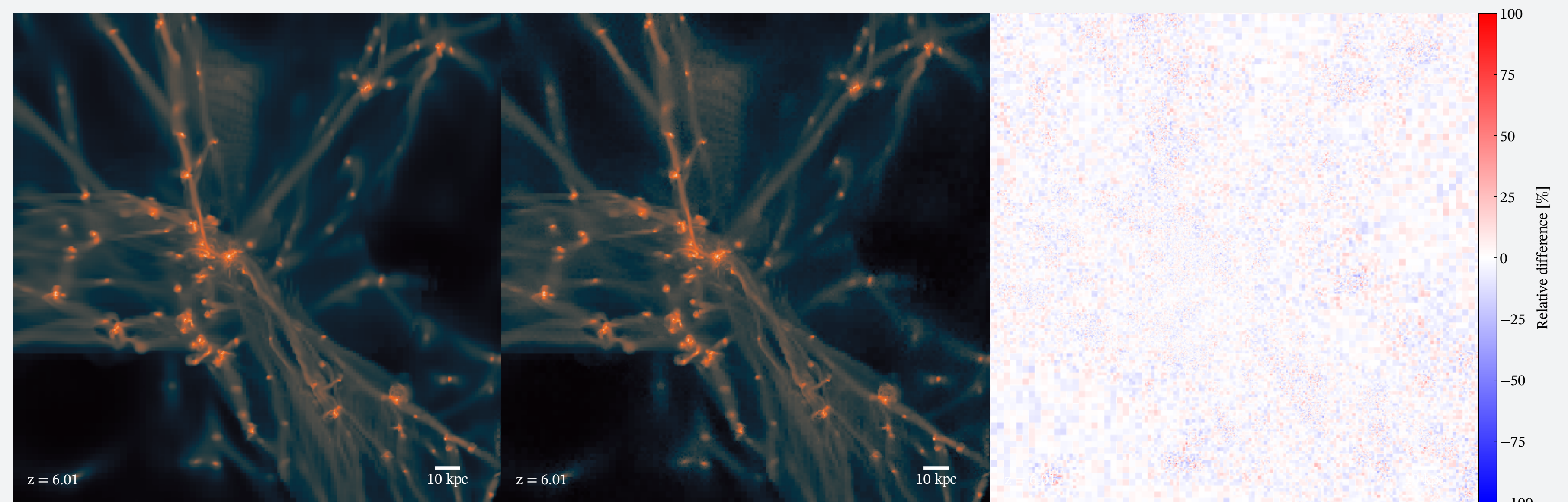
3D Sedov blast. Solid: gas radial profile at different times. Symbols: tracer density profile with error bars at $\pm 10\sigma$.

STARS



Star formation: gas tracers are attached onto stars at star formation using $p_{i\star} = M_\star/M_i$.
Supernova feedback: star tracers are released with probability $p_{\star i} = \eta$. Tracers are released in the up-to-48 neighboring cells.

MONTE CARLO TRACERS – RESULTS



The gas tracer distribution (center) matches the gas distribution (left) at percent level, **regardless of the geometry of the flow** (right, same scale as for the velocity tracers). The distribution of the number of tracers per cell is given by a Poisson law

$$p(N_{\text{tracer}} = k) \sim \text{Poisson}(\lambda = M_{\text{cell}}/m_{\text{tracer}}) \equiv e^{-\lambda} \frac{\lambda^k}{k!}. \quad (2)$$

The distribution of number of tracers per star is well approximated by a Poisson law with parameter $\lambda = M_\star/m_{\text{tracer}}$.

The tracer particles are correctly interfaced with gas, star formation, SN feedback and AGN feedback.